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# A precontact binary and a shallow contact binary are in the same field 

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#### Abstract

The period changes of two close binaries, V1107 Cas and AX Cas, which are in the same field, were investigated. Their periods both show a long-term decrease. After further analysis, we found that the periods have their respective cyclic oscillations ( $T_{3}=6.74 \pm$ 0.24 yr for V 1107 Cas and $T_{3}=13.8 \pm 0.3 \mathrm{yr}$ for AX Cas), which are possibly caused by a third body due to the light-time effect. We also obtained the complete $V R_{c} I_{\mathrm{c}}$ light curves for V1107 Cas and analyzed them with the 2010 version of the Wilson-Devinney code. The photometric results reveal that V1107 Cas is a W-type shallow contact ( $15.2 \% \pm 1.8 \%$ ) binary, with a mass-ratio of $1.797 \pm 0.006$. The period variation and photometric solution suggest that V1107 Cas is a newly formed contact binary system. Moreover, we estimated the fundamental parameters for V1107 Cas. They are: $M_{1}=0.39 \pm 0.01 M_{\odot}, M_{2}=0.70 \pm$ $0.03 M_{\odot}, R_{1}=0.52 \pm 0.10 R_{\odot}, R_{2}=0.68 \pm 0.12 R_{\odot}, L_{1}=0.178 \pm 0.108 L_{\odot}$, and $L_{2}=0.196$ $\pm 0.116 L_{\odot}$. Then, based on the coplane assumption, we deduced the masses of possible third bodies to be $M_{3}=0.091 \pm 0.019 M_{\odot}$ for V1107 Cas and $M_{3}=0.325 \pm 0.029 M_{\odot}$ for AXCas. Finally, we inferred the evolutional stage of AXCas, and believe that it is a precontact binary. Thus, the precontact binary AXCas and the shallow contact binary V1107 Cas have adjoining evolutional stages.


Key words: binaries: close — binaries: eclipsing — stars: evolution — stars: individuals (V1107 Cassiopeiae; AX Cassiopeiae)

## 1 Introduction

Mass transfer is a common and very important process for close binaries, which drives the evolution of the systems, and had been used to explain the famous Algol paradox, as well many other interesting astrophysical processes. Understanding the mass transfer that occurs in near contact binaries or in marginal contact binaries is key to understanding how a semi-detached binary evolves into a contact binary
(i.e., BL And and GW Tau, Zhu \& Qian 2006; AS Ser, Zhu et al. 2008; DD Com, Zhu et al. 2010). Some such masstransfer processes were affected by third bodies or by magnetic activeties (i.e., ZZ Eri, D. R. Faulkner et al. 2014; ${ }^{1}$ GSC 1537-1557, AHTau, and V508 Oph, Xiang et al. 2015a, 2015b, 2015c). The absolute physical parameters
${ }^{1}$ Faulkner, D. R., Clark, J., Samec, R. G., Hill, R. L., Kring, J., Flaaten, D., \& Van Hamme, W. V. 2014, Abstract, AAS Meeting \#223, id.155.11.
that combine the orbital elements are a great help for dig－ ging out evolutional information concerning a close binary， so we monitored the times of the minima for V1107 Cas， and observed its multiple color light curves，while intending to find the evolutional stage of V1107 Cas．During the data processing，we realized that the other close binary，AX Cas， had been observed in the same field，by chance．AX Cas had not been included in our observational schedule because it is not an EW but EB－type binary．However，after having investigated some of the literature（i．e．，AK CMi，Samec et al．1998；V2421 Cyg，Samec et al．2014），we became aware that AX Cas may be a precontact binary．It is very interesting that a precontact and a contact binary are in the same field．

Until and after V1107 Cas（GSC 04030－02020）had been found as a W UMa variable（Hoffman et al．2009）， there were only 10 papers written about it；none of them as an individual study．From the 2 Micron All－Sky Survey （2MASS）catalogue（R．M．Cutri et al．2003），${ }^{2}$ we found its $J H K$ magnitudes to be $J=11.250 \pm 0.023, H=10.806$ $\pm 0.024$ ，and $K=10.717 \pm 0.018$ ．And according to the NOMAD Catalog（N．Zacharias et al．2005），${ }^{3}$ we found its $B V R$ magnitudes to be $B=13.790, V=13.310$ ，and $R$ $=12.770$ ．The yielded color indexes by these magnitudes were used to estimate the effective temperature of star 1 in our photometric solution．As for the other close binary， AX Cas，it was discovered as early as 1931 by Beljawsky （1931）．We found its $J H K$ magnitudes to be $J=10.991$ $\pm 0.022, H=10.655 \pm 0.019$ ，and $K=10.532 \pm 0.016$ （R．M．Cutri et al．2003）， 2 and its $B V R$ magnitudes to be $B=13.010, V=12.640$ ，and $R=12.180$（N．Zacharias et al．2005）．${ }^{3}$ We computed the masses，radii，and absolute thermal magnitudes of its components as $M_{1}=2.0 M_{\odot}, M_{2}$ $=0.45 M_{\odot}, R_{1}=2.10 R_{\odot}, R_{2}=1.00 R_{\odot}, M_{\mathrm{bol} 1}=1.35 \mathrm{mag}$ ， $M_{\mathrm{bol} 2}=4.80 \mathrm{mag}$（M．A．Svechnikov \＆Eh．F．Kuznetsova 1990）．${ }^{4}$ However，the color indexes did not agree with the less－massive component．We discuss this in section 5.

## 2 New CCD imaging photometry

We obtained complete $V R_{\mathrm{c}} I_{\mathrm{c}}$ light curves of V1107 Cas on 2013 August 18 and 25 ，using the $1024 \times 1024$ PI1024 BFT

[^0]camera attached to the 85 cm telescope at the Xinglong Sta－ tion of the National Astronomical Observatories of the Chi－ nese Academy of Sciences．The observational system has a standard Johnson－Cousins－Bessel multicolor CCD photo－ metric system built on the primary focus（Zhou et al．2009）， generating an effective field of view equal to $16!5 \times 16!5$ ． We also obtained several times of minima of V1107 Cas， using the DW436 $2048 \times 2048$ CCD photometric system attached to the 1 m reflecting telescope and the DV436 2048 $\times 2048$ CCD photometric system attached to the 60 cm reflecting telescope at the Yunnan Observatories（YNOs） in China．Both filter systems were the Johnson－Cousins－ Bessel system．The effective field of view of the 1 m tele－ scope was $7!3 \times 7!3$ ，while that of the 60 cm telescope was $12^{\prime} \times 12^{\prime}$ at the Cassegrain focus．These times of minima of V1107 Cas are listed in table 1，where＂ 1 m ＂and＂ 60 cm ＂ refer to the times of minima monitored with the 1 m and 60 cm telescopes at the YNOs，respectively；＂ 85 cm ＂ refers to those monitored with the 85 cm telescope at the Xinglong Station．2MASS J01233287＋6135269（star 1） and 2MASS J01232239＋6131552（star 2）were used for photometric comparisons．The PHOT／IRAF aperture pho－ tometry package was exercised so as to reduce the images， using the standard procedure of flat－fielding，except for dark corrections，because the dark value of the images was very insignificant．The corresponding light curves of V1107 Cas are shown in figure 1 ，where the phases were calculated with equation（6）．On the other hand，the light levels of the curves taken on two nights match very well，suggesting little nightly variation．On the other hand，the magnitude differ－ ences between the two comparison stars remained nearly constant，vindicating the assumption that the comparisons are nonvariable．A quadratic polynomial model was used to determine the times of minima for V1107 Cas by the least－squares method．Our new times of minima are listed in table 1．Because the other close binary AX Cas is close to V1107 Cas in visual，it was observed at the same time in most of our frames．Although we did not obtain complete light curves of AX Cas，we obtained some times of minima． They are also listed in table 1.

## 3 Photometric solutions of V1107 Cas

Multicolor light curves of V1107 Cas were analyzed with the 2010 Version of the Wilson－Devinney（W－D）code （Wilson \＆Devinney 1971；Wilson 1979，1990，1994， 2008；Van Hamme \＆Wilson 2007）．The effective temper－ ature of star 1 was estimated from the colors of V1107 Cas mentioned in section 1 using the program of Worthey and Lee（2011）．It should range between 4900 K and 5500 K ． In the solution，we took the temperature as 5200 K ，and then used the temperature range to limit its uncertainty．

Table 1. New times of light minimum for AX Cas and V1107 Cas.

| JD (Hel.) (d) | Error (d) | Min. | Filter | Telescope |
| :--- | :---: | :---: | :---: | :---: |
|  |  | -AX Cas- |  |  |
| 2456214.22009 | 0.00011 | sec. | No | 1 m |
| 2456220.22436 | 0.00038 | sec. | $V$ | 1 m |
| 2456220.22450 | 0.00034 | prim. | $R_{\mathrm{c}}$ | 1 m |
| 2456220.22467 | 0.00033 | prim. | $I_{\mathrm{c}}$ | 1 m |
| 2456220.22476 | 0.00041 | sec. | $B$ | 1 m |
|  |  | - V1107 Cas- |  |  |
|  |  |  |  |  |
| 2456209.16515 | 0.00011 | sec. | No | 60 cm |
| 2456214.21991 | 0.00012 | prim. | No | 1 m |
| 2456220.23732 | 0.00021 | prim. | $V$ | 1 m |
| 2456220.23741 | 0.00020 | prim. | $R_{\mathrm{c}}$ | 1 m |
| 2456220.23753 | 0.00019 | prim. | $I_{\mathrm{c}}$ | 1 m |
| 2456220.23755 | 0.00025 | prim. | $B$ | 1 m |
| 2456220.37551 | 0.00045 | sec. | $B$ | 1 m |
| 2456220.37568 | 0.00026 | sec. | $R_{\mathrm{c}}$ | 1 m |
| 2456220.37570 | 0.00031 | sec. | $V$ | 1 m |
| 2456220.37583 | 0.00029 | sec. | $I_{\mathrm{c}}$ | 1 m |
| 2456523.31517 | 0.00015 | sec. | $I_{\mathrm{c}}$ | 85 cm |
| 2456523.31517 | 0.00011 | sec. | $V$ | 85 cm |
| 2456523.31522 | 0.00013 | sec. | $R_{\mathrm{c}}$ | 85 cm |
| 2456530.28646 | 0.00009 | prim. | $I_{\mathrm{c}}$ | 85 cm |
| 2456530.28673 | 0.00009 | prim. | $V$ | 85 cm |
| 2456530.28691 | 0.00008 | prim. | $R_{\mathrm{c}}$ | 85 cm |



Fig. 1. $V R_{\mathrm{c}} I_{\mathrm{c}}$-band light curves of V 1107 Cas taken on 2012 August 18 and 25. The light levels of the curves taken on the two nights match very well, suggesting little nightly variation. The two-day light curves are phased up smoothly, suggesting little nightly fluctuations.

We thus have $T_{1}=5200 \pm 300 \mathrm{~K}$. We then used the $q$-search method to obtain an initial input value of the mass ratio, $q$. The details are: we fixed the mass ratio $q$ at a series of values of $0.1,0.2,0.3$, etc.; we then fitted the light curves with the W-D code for each $q$ value, obtaining a series of corresponding fitting residuals, illustrated in figure 2, and selected the final $q$ value with minimal residuals. The mass


Fig. 2. Relationship between the mass ratio, $q$, and the fitting residual.
ratio should thus be between 1.6 and 3.2. While obtaining the solution, the bolometric albedo $A_{1}=A_{2}=0.5$ (Rucinski 1969) and the values of the gravity-darkening coefficient $g_{1}$ $=g_{2}=0.32$ (Lucy 1967) were used, which correspond to the common convective envelope of both components. According to Claret and Gimenez (1990), the square-root limb-darkening coefficients were used. We adjusted the mass ratio, $q$; the mean temperature of star $2, T_{2}$; the monochromatic luminosity of star 1 and the dimensionless

Table 2. Photometric solutions for V1107 Cas.

| Parameters | Without $l_{3}$ $V R_{\mathrm{c}} I_{\mathrm{c}}$ | Errors | With $l_{3}$ $V R_{\mathrm{c}} I_{\mathrm{c}}$ | Errors |
| :---: | :---: | :---: | :---: | :---: |
| $g_{1}=g_{2}$ | 0.32 | fixed | 0.32 | fixed |
| $A_{1}=A_{2}$ | 0.50 | fixed | 0.50 | fixed |
| $x_{1 \text { bolo }}=x_{2 \mathrm{bolo}}, y_{1 \mathrm{bolo}}=y_{2 \mathrm{bolo}}$ | 0.271, 0.433 | fixed | 0.271, 0.433 | fixed |
| $x_{1 V}=x_{2 V}, y_{1 V}=y_{2 V}$ | 0.468, 0.373 | fixed | 0.468, 0.373 | fixed |
| $x_{1 R_{\mathrm{c}}}=x_{2 R_{\mathrm{c}}}, y_{1 R_{\mathrm{c}}}=y_{2 R_{\mathrm{c}}}$ | 0.307, 0.483 | fixed | 0.307, 0.483 | fixed |
| $x_{1 I_{\mathrm{c}}}=x_{2 I_{\mathrm{c}}}, y_{1 I_{\mathrm{c}}}=y_{2 I I_{\mathrm{c}}}$ | 0.199, 0.508 | fixed | 0.199, 0.508 | fixed |
| $T_{1}(\mathrm{~K})$ | 5200 | $\pm 300$,fixed | 5200 | $\pm 300$,fixed |
| $T_{2}(\mathrm{~K})$ | 4657 | $\pm 306$ | 4657 | $\pm 306$ |
| $q=M_{2} / M_{1}$ | 1.808 | $\pm 0.005$ | 1.797 | $\pm 0.006$ |
| $i\left({ }^{\circ}\right)$ | 72.0 | $\pm 0.1$ | 72.3 | $\pm 0.5$ |
| $\Omega_{1}=\Omega_{2}$ | 4.8898 | $\pm 0.0095$ | 4.8721 | $\pm 0.0106$ |
| $\Omega_{\text {in }}$ | 4.9773 | - | 4.9615 | - |
| $\Omega_{\text {out }}$ | 4.3866 | - | 4.3713 | - |
| $L_{1} /\left(L_{1}+L_{2}\right)(V)$ | 0.5252 | $\pm 0.0020$ | 0.5266 | $\pm 0.0114$ |
| $L_{1} /\left(L_{1}+L_{2}\right)\left(R_{\mathrm{c}}\right)$ | 0.4945 | $\pm 0.0016$ | 0.4959 | $\pm 0.0106$ |
| $L_{1} /\left(L_{1}+L_{2}\right)\left(I_{\mathrm{c}}\right)$ | 0.4757 | $\pm 0.0013$ | 0.4772 | $\pm 0.0102$ |
| $l_{3} /\left(L_{1}+L_{2}+l_{3}\right)(V)$ | - | - | 0.0011 | $\pm 0.0114$ |
| $l_{3} /\left(L_{1}+L_{2}+l_{3}\right)\left(R_{\mathrm{c}}\right)$ | - | - | 0.0006 | $\pm 0.0106$ |
| $l_{3} /\left(L_{1}+L_{2}+l_{3}\right)\left(I_{\mathrm{c}}\right)$ | - | - | 0.0016 | $\pm 0.0102$ |
| $r_{1}$ (pole) | 0.3159 | $\pm 0.0010$ | 0.3166 | $\pm 0.0012$ |
| $r_{1}$ (side) | 0.3312 | $\pm 0.0012$ | 0.3319 | $\pm 0.0014$ |
| $r_{1}$ (back) | 0.3688 | $\pm 0.0020$ | 0.3697 | $\pm 0.0023$ |
| $r_{2}$ (pole) | 0.4137 | $\pm 0.0009$ | 0.4134 | $\pm 0.0010$ |
| $r_{2}$ (side) | 0.4402 | $\pm 0.0012$ | 0.4399 | $\pm 0.0013$ |
| $r_{2}$ (back) | 0.4721 | $\pm 0.0016$ | 0.4719 | $\pm 0.0018$ |
| $f(\%)$ | 14.8 | $\pm 1.6$ | 15.2 | $\pm 1.8$ |

potential of star 1. We started the differential corrections (DC) program at model 2. After thousands of iterations, it converged to model 3 ( $\Omega_{1}=\Omega_{2}$, mode 3 for contact configuration). The photometric solution is listed in table 2 and the theoretical light curves computed with those photometric elements are plotted in figure 3. We also ran the program with $l_{3}$ light, because the period investigation of V1107 Cas, described in next section, implied that it has a third companion. The outcomes (table 2) showed that the third companion has no more than a $0.16 \%$ luminosity contribution to the whole system. Its details are given in section 5 .

According to the photometric solution, we concluded that V1107 Cas is a moderate mass-ratio ( $1.797 \pm 0.006$ ), shallow contact ( $15.2 \% \pm 1.8 \%$ ) binary, where "shallow" is defined when the contact degree is less than $20 \%$. We assumed that primary component of V1107 Cas is a normal main-sequence star. Accordingly, we found its temperature to be $T=4657 \pm 306 \mathrm{~K}$, and estimated its mass at $M=$ $0.70 \pm 0.03$. We then estimated the physical parameters of each component by using of the light-curve (LC) program, W-D code (Wilson \& Devinney 1971; Wilson 1979, 1990, 1994, 2008; Van Hamme \& Wilson 2007). The results are


Fig. 3. Observed and theoretical light curves in the $V, R_{\mathrm{c}}$, and $I_{\mathrm{c}}$ bands of V1107 Cas, with $I_{3}$ light.
$M_{1}=0.39 \pm 0.01 M_{\odot}, M_{2}=0.70 \pm 0.03 M_{\odot}, R_{1}=0.52 \pm$ $0.10 R_{\odot}, R_{2}=0.68 \pm 0.12 R_{\odot}, L_{1}=0.178 \pm 0.108 L_{\odot}$, and $L_{2}=0.196 \pm 0.116 L_{\odot}$. The estimated uncertainties of the absolute parameters mainly result from the uncertainties of the temperatures of the components.

Table 3. CCD Times of light minima for V1107 Cas.

| JD (Hel.) | Filters | Min. | Epoch | $(\mathrm{O}-\mathrm{C})(\mathrm{d})$ | $(\mathrm{O}-\mathrm{C})_{1}(\mathrm{~d})$ | Residual (d) | Reference |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2453215.4410 |  | prim. | -12124 | -0.00304 | -0.00062 | -0.00058 | IBVS No. 5731 |
| 2453215.5786 |  | sec. | -12123.5 | -0.00214 | 0.00027 | 0.00031 | IBVS No. 5731 |
| 2453254.4031 |  | sec. | -11981.5 | -0.00211 | 0.00018 | 0.00038 | IBVS No. 5731 |
| 2453637.4522 |  | sec. | -10580.5 | -0.00290 | -0.00165 | -0.00020 | IBVS No. 5731 |
| 2454085.3015 |  | sec. | -8942.5 | -0.00208 | -0.00176 | -0.00050 | IBVS No. 5761 |
| 2454092.2753 |  | prim. | -8917 | -0.00028 | 0.00001 | 0.00126 | IBVS No. 5874 |
| 2454092.4111 |  | sec. | -8916.5 | -0.00118 | -0.00088 | 0.00036 | IBVS No. 5874 |
| 2454092.5478 |  | prim. | -8916 | -0.00119 | -0.00089 | 0.00035 | IBVS No. 5874 |
| 2454092.6810 |  | sec. | -8915.5 | -0.00469 | -0.00439 | -0.00315 | IBVS No. 5874 |
| 2454308.4086 |  | sec. | -8126.5 | 0.00102 | 0.00097 | 0.00149 | IBVS No. 5874 |
| 2454367.3284 |  | prim. | -7911 | 0.00058 | 0.00045 | 0.00074 | IBVS No. 5874 |
| 2454367.4657 |  | sec. | -7910.5 | 0.00118 | 0.00105 | 0.00133 | IBVS No. 5874 |
| 2454367.6015 |  | prim. | -7910 | 0.00027 | 0.00014 | 0.00042 | IBVS No. 5874 |
| 2454388.3799 |  | prim. | -7834 | -0.00062 | -0.00077 | -0.00057 | IBVS No. 5874 |
| 2454388.5175 |  | sec. | -7833.5 | 0.00027 | 0.00011 | 0.00031 | IBVS No. 5874 |
| 2454388.6539 |  | prim. | -7833 | -0.00003 | -0.00018 | 0.00000 | IBVS No. 5874 |
| 2454390.2941 |  | prim. | -7827 | -0.00030 | -0.00046 | -0.00027 | IBVS No. 6070 |
| 2454673.5509 |  | prim. | -6791 | 0.00190 | 0.00143 | 0.00051 | IBVS No. 5918 |
| 2454704.8561 |  | sec. | -6676.5 | 0.00146 | 0.00096 | -0.00006 | IBVS No. 5875 |
| 2454776.3525 |  | prim. | -6415 | 0.00068 | 0.00013 | -0.00110 | IBVS No. 5918 |
| 2454776.4901 |  | sec. | -6414.5 | 0.00157 | 0.00102 | -0.00021 | IBVS No. 5918 |
| 2454835.2734 |  | sec. | -6199.5 | 0.00134 | 0.00075 | -0.00062 | IBVS No. 5918 |
| 2454847.3039 |  | sec. | -6155.5 | 0.00173 | 0.00113 | -0.00027 | IBVS No. 5918 |
| 2454847.4412 |  | prim. | -6155 | 0.00232 | 0.00172 | 0.00031 | IBVS No. 5918 |
| 2454847.5771 |  | sec. | -6154.5 | 0.00151 | 0.00092 | -0.00048 | IBVS No. 5918 |
| 2455081.3444 |  | sec. | -5299.5 | 0.00175 | 0.00106 | -0.00055 | IBVS No. 5941 |
| 2455081.4805 |  | prim. | -5299 | 0.00115 | 0.00045 | -0.00116 | IBVS No. 5941 |
| 2455081.6183 |  | sec. | -5298.5 | 0.00224 | 0.00154 | -0.00007 | IBVS No. 5941 |
| 2455154.3459 |  | sec. | -5032.5 | 0.00231 | 0.00160 | 0.00003 | IBVS No. 5941 |
| 2455154.4840 |  | prim. | -5032 | 0.00370 | 0.00299 | 0.00143 | IBVS No. 5941 |
| 2455154.6182 |  | sec. | -5031.5 | 0.00120 | 0.00049 | -0.00107 | IBVS No. 5941 |
| 2455374.4429 |  | sec. | -4227.5 | 0.00284 | 0.00213 | 0.00103 | IBVS No. 5984 |
| 2455409.4380 |  | sec. | -4099.5 | 0.00123 | 0.00054 | -0.00045 | IBVS No. 5984 |
| 2455479.2941 |  | prim. | -3844 | 0.00062 | -0.00004 | -0.00079 | IBVS No. 5984 |
| 2455479.4324 |  | sec. | -3843.5 | 0.00222 | 0.00154 | 0.00080 | IBVS No. 5984 |
| 2455479.5689 |  | prim. | -3843 | 0.00201 | 0.00134 | 0.00059 | IBVS No. 5984 |
| 2455482.3017 | R | prim. | -3833 | 0.00069 | 0.00002 | -0.00070 | B.R.N.O. No. 37 |
| 2455482.4394 | $R$ | sec. | -3832.5 | 0.00169 | 0.00102 | 0.00028 | B.R.N.O. No. 37 |
| 2455491.3259 |  | prim. | -3800 | 0.00231 | 0.00164 | 0.00093 | IBVS No. 5984 |
| 2455491.4624 |  | sec. | -3799.5 | 0.00210 | 0.00143 | 0.00073 | IBVS No. 5984 |
| 2455491.5976 |  | prim. | -3799 | 0.00059 | -0.00006 | -0.00077 | IBVS No. 5984 |
| 2455514.4290 |  | sec. | -3715.5 | 0.00211 | 0.00145 | 0.00084 | IBVS No. 5984 |
| 2455514.5662 |  | prim. | -3715 | 0.00260 | 0.00195 | 0.00133 | IBVS No. 5984 |
| 2455792.3489 | R | prim. | -2699 | -0.00105 | -0.00153 | -0.00101 | B.R.N.O. No. 38 |
| 2455792.4862 | R | sec. | -2698.5 | -0.00046 | -0.00093 | -0.00042 | B.R.N.O. No. 38 |
| 2455794.4008 | R | sec. | -2691.5 | 0.00026 | -0.00021 | 0.00030 | B.R.N.O. No. 38 |
| 2455794.5370 | R | prim. | -2691 | -0.00025 | -0.00072 | -0.00020 | B.R.N.O. No. 38 |
| 2455796.4504 | $R$ | prim. | -2684 | -0.00073 | -0.00120 | -0.00067 | B.R.N.O. No. 38 |
| 2455796.5871 | R | sec. | -2683.5 | -0.00073 | -0.00120 | -0.00068 | B.R.N.O. No. 38 |
| 2455797.4080 | R | sec. | -2680.5 | -0.00007 | -0.00054 | -0.00001 | B.R.N.O. No. 38 |
| 2455797.5435 | $R$ | prim. | -2680 | -0.00128 | -0.00175 | -0.00121 | B.R.N.O. No. 38 |
| 2455799.3220 | $R$ | sec. | -2673.5 | 0.00005 | -0.00042 | 0.00011 | B.R.N.O. No. 38 |
| 2455799.4579 | R | prim. | -2673 | -0.00076 | -0.00123 | -0.00069 | B.R.N.O. No. 38 |
| 2455799.5947 | R | sec. | -2672.5 | -0.00066 | -0.00113 | -0.00059 | B.R.N.O. No. 38 |
| 2455802.3292 | $R$ | sec. | -2662.5 | -0.00028 | -0.00075 | -0.00020 | B.R.N.O. No. 38 |

Table 3 (Continued)

| JD (Hel.) | Filters | Min. | Epoch | $(\mathrm{O}-\mathrm{C})(\mathrm{d})$ | $(\mathrm{O}-\mathrm{C})_{1}(\mathrm{~d})$ | Residual (d) | Reference |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2455802.4650 | $R$ | prim. | -2662 | -0.00119 | $-0.00165$ | -0.00110 | B.R.N.O. No. 38 |
| 2455805.3369 | $R$ | sec. | -2651.5 | -0.00011 | -0.00057 | -0.00001 | B.R.N.O. No. 38 |
| 2455805.4725 | $R$ | prim. | -2651 | -0.00122 | -0.00168 | -0.00112 | B.R.N.O. No. 38 |
| 2455805.6098 | R | sec. | -2650.5 | -0.00062 | -0.00109 | -0.00052 | B.R.N.O. No. 38 |
| 2455815.3159 | $R$ | prim. | -2615 | -0.00064 | -0.00109 | -0.00049 | B.R.N.O. No. 38 |
| 2455815.4536 | $R$ | sec. | -2614.5 | 0.00035 | -0.00010 | 0.00049 | B.R.N.O. No. 38 |
| 2455830.3531 | $R$ | prim. | -2560 | -0.00109 | -0.00153 | -0.00087 | B.R.N.O. No. 38 |
| 2455830.4924 | R | sec. | -2559.5 | 0.00151 | 0.00106 | 0.00172 | B.R.N.O. No. 38 |
| 2455831.3118 | $R$ | sec. | -2556.5 | 0.00067 | 0.00022 | 0.00088 | B.R.N.O. No. 38 |
| 2455831.4473 | $R$ | prim. | -2556 | -0.00053 | -0.00097 | -0.00031 | B.R.N.O. No. 38 |
| 2455835.4135 |  | sec. | -2541.5 | 0.00119 | 0.00075 | 0.00143 | IBVS No. 6070 |
| 2455838.2831 | $R$ | prim. | -2531 | -0.00003 | -0.00046 | 0.00022 | B.R.N.O. No. 38 |
| 2455838.4203 | $R$ | sec. | -2530.5 | 0.00047 | 0.00002 | 0.00071 | B.R.N.O. No. 38 |
| 2455852.2271 | R | prim. | -2480 | -0.00003 | -0.00045 | 0.00028 | B.R.N.O. No. 38 |
| 2455858.2432 |  | prim. | -2458 | 0.00101 | 0.00059 | 0.00135 | IBVS No. 6026 |
| 2455858.3804 |  | sec. | -2457.5 | 0.00151 | 0.00108 | 0.00184 | IBVS No. 6026 |
| 2455858.5156 |  | prim. | -2457 | 0.00000 | -0.00042 | 0.00034 | IBVS No. 6026 |
| 2455875.3307 | $R$ | sec. | -2395.5 | 0.00028 | -0.00012 | 0.00069 | B.R.N.O. No. 38 |
| 2455893.2382 |  | prim. | -2330 | -0.00069 | -0.00108 | -0.00019 | IBVS No. 6026 |
| 2455905.6777 |  | sec. | -2284.5 | -0.00143 | -0.00180 | -0.00087 | IBVS No. 6011 |
| 2456154.3465 | $R$ | prim. | -1375 | -0.00063 | -0.00070 | 0.00083 | B.R.N.O. No. 38 |
| 2456154.4831 | R | sec. | -1374.5 | -0.00074 | -0.00081 | 0.00072 | B.R.N.O. No. 38 |
| 2456179.3608 |  | sec. | -1283.5 | -0.00351 | -0.00355 | -0.00198 | IBVS No. 6070 |
| 2456202.8746 | V | sec. | -1197.5 | -0.00312 | -0.00312 | -0.00153 | IBVS No. 6042 |
| 2456209.16515 | $N$ | sec. | -1174.5 | -0.00104 | -0.00103 | 0.00055 | The paper |
| 2456214.21991 | $N$ | prim. | -1156 | -0.00440 | -0.00439 | -0.00278 | The paper |
| 2456220.23745 | $B V R_{\text {c }} I_{\text {c }}$ | prim. | -1134 | -0.00192 | -0.00190 | -0.00029 | The paper |
| 2456220.3755 |  | sec. | -1133.5 | -0.00057 | -0.00055 | 0.00104 | IBVS No. 6070 |
| 2456220.37568 | $B V R_{\text {c }} I_{\text {c }}$ | sec. | -1133.5 | -0.00039 | -0.00037 | 0.00122 | The paper |
| 2456222.2889 |  | sec. | -1126.5 | -0.00106 | -0.00103 | 0.00057 | IBVS No. 6070 |
| 2456222.4247 |  | prim. | -1126 | -0.00196 | -0.00194 | -0.00033 | IBVS No. 6070 |
| 2456523.31519 | $V R_{\text {c }} I_{\text {c }}$ | sec. | -25.5 | -0.00112 | -0.00058 | 0.00074 | The paper |
| 2456530.2867 | $V R_{\mathrm{c}} I_{\text {c }}$ | prim. | 0 | -0.00161 | -0.00106 | 0.00024 | The paper |

## 4 Orbital period analysis

### 4.1 V1107Cas

Data of times of minima for this contact binary from 2004 to 2013 were collected. All of these 88 times of minima are CCD data, and are listed in table 3, which were used to correct the period with the ephemeris (HJD) 2456530.2867 (To) $+0 .{ }^{\mathrm{d}} 273406\left(P_{0}\right) \times E$, where $T_{0}$ is one of our minimum value, and $P_{0}$ was obtained from GCVS (N. N. Samus et al. 2009). ${ }^{5}$ The new linear ephemeris is

$$
\begin{align*}
\operatorname{Min} . I(H J D)= & 2456530.2883( \pm 0.0003) \\
& +0.27341177( \pm 0.00000006) \times E \tag{1}
\end{align*}
$$

Because the $(O-C)$ diagram derived from this formula shows a long-term decreasing tendency, we at first used a parabola term to fit this $(O-C)$ diagram. However, we found that the residuals of this fitting still showed a cyclical variation. Finally, we used a quadratic change superimposed upon a cyclical change to fit the $(O-C)$ diagram, yielding a fitting of

Min. $I(H J D)=2456530.2878( \pm 0.0006)$

$$
\begin{align*}
& +0.27341123( \pm 0.00000024) \times E \\
& -5.69( \pm 2.06) \times 10^{-11} \times E^{2} \\
& +0.0016( \pm 0.0001) \sin (0.040 E-53.6) \tag{2}
\end{align*}
$$

The resultant $(O-C)$ and $(O-C)_{1}$ diagrams are also shown in figures 4 and 5 , as the solid line represents. That

[^1]

Fig. 4. $(O-C)$ diagram of V1107 Cas formed by all available measurements. The $(O-C)$ values were computed by using a newly determined linear ephemeris, equation (1). The solid circles refer to the primary minima, and the open ones to the secondary minima; the solid line represents a long-term decrease variation, equation (2). The residuals derived from equation (2) are displayed in the lower panel.


Fig. 5. $(O-C)_{1}$ diagram of V 1107 Cas. The symbols are the same as those in figure 4.
fitting gives a long-term decrease $[d P / d t=-1.52( \pm 0.55) \times$ $\left.10^{-7} \mathrm{~d} \mathrm{yr}^{-1}\right]$ with a clear period oscillation $\left[A_{3}=0.0016\right.$, $T_{3}=6.74 \mathrm{yr} ; T_{3}$ was computed by equation (9)]. The residuals are plotted at the bottom of figure 4 .

The period decrease should be due to mass transfer from the more-massive component to the less-massive one. If this mass transfer is conservative, we can estimate the masstransfer rate by using a well-known equation (Tout \& Hall 1991), with the absolute parameters derived by the present
paper. The formula is

$$
\begin{equation*}
\frac{\dot{P}}{P}=3 \frac{\dot{M_{2}}}{M_{2}}\left(\frac{M_{2}}{M_{1}}-1\right), \tag{3}
\end{equation*}
$$

so the mass-transfer rate is estimated to be, $d M_{2} / d t=$ $-1.61( \pm 0.63) \times 10^{-7} M_{\odot} \mathrm{yr}^{-1}$. The negative sign implies that the more-massive component is losing matter. The time-scale of mass transfer is $\tau \sim M_{2} /\left|\dot{M}_{2}\right| \sim 4.3( \pm 1.7) \times$ $10^{6} \mathrm{yr}$, which is far greater than the thermal time scale $\left(t_{\mathrm{th}}=2 \times 10^{7} M^{2} R^{-1} L^{-1} \mathrm{yr}\right.$, where $M, R$, and $L$ are all in solar units) of the more-massive component [ $t_{\mathrm{th}}=$ $\left.7.4( \pm 4.6) \times 10^{7} \mathrm{yr}\right]$. The mass-transfer rate is so great that the less-massive component accepting the matter cannot redistribute it in time, so that the temperature of the less-massive component rises. That is why the temperature difference between the two components is as high as 543 K .

As shown in figures 4 and 5, both the primary and the secondary times of light minimum light follow the same general trend of an $(O-C)$ variation, indicating that the ( $O-C$ ) oscillation cannot be explained as apsidal motion. Because V1107Cas contains later-type components, its alternate period change can be interpreted in term of the mechanism of magnetic activity (e.g., Applegate 1992; Lanza et al. 1998). However, according exhaustive statistical work, Liao and Qian (2010) pointed out that the

Table 4. Masses and orbital radii of the assumed third bodies in AXCas and V1107 Cas.

| Parameters | $\begin{gathered} \text { AX Cas } \\ M_{1}=2.0 M_{\odot} \end{gathered}$ | $\begin{gathered} \text { AX Cas } \\ M_{1}=0.9 M_{\odot} \end{gathered}$ | V1107 Cas | Units |
| :---: | :---: | :---: | :---: | :---: |
| $A_{3}$ | 0.0084( $\pm 0.0044)$ | 0.0084( $\pm 0.0044)$ | 0.0016( $\pm 0.0001)$ | d |
| $T_{3}$ | $13.8( \pm 0.3)$ | $13.8( \pm 0.3)$ | $6.74( \pm 0.24)$ | yr |
| $e^{\prime}$ | 0.2533( $\pm 0.1167)$ | 0.2533 ( $\pm 0.1167)$ | 0 (assumed) | - |
| $w^{\prime}$ | 87.8(土4.6) | 87.8(土4.6) | - | $\bigcirc$ |
| $a_{12}^{\prime} \sin i^{\prime}$ | $1.45( \pm 0.19)$ | $1.45( \pm 0.19)$ | $0.28( \pm 0.10)$ | au |
| $f(m)$ | $1.61( \pm 0.94) \times 10^{-2}$ | $1.61( \pm 0.94) \times 10^{-2}$ | $4.68( \pm 0.9) \times 10^{-4}$ | $M_{\odot}$ |
| $m_{3}\left(i^{\prime}=90^{\circ}\right)$ | 0.509 ( $\pm 0.028)$ | $0.319( \pm 0.028)$ | 0.086( $\pm 0.018)$ | $M_{\odot}$ |
| $m_{3}\left(i^{\prime}=70^{\circ}\right)$ | 0.546( $\pm 0.029)$ | 0.344( $\pm 0.029)$ | 0.092( $\pm 0.019)$ | $M_{\odot}$ |
| $m_{3}\left(i^{\prime}=50^{\circ}\right)$ | 0.692( $\pm 0.046)$ | $0.440( \pm 0.046)$ | 0.114( $\pm 0.066)$ | $M_{\odot}$ |
| $m_{3}\left(i^{\prime}=30^{\circ}\right)$ | $1.169( \pm 0.211)$ | $0.766( \pm 0.211)$ | 0.182( $\pm 0.011)$ | $M_{\odot}$ |
| $m_{3}\left(i^{\prime}=10^{\circ}\right)$ | 5.975 ( $\pm 0.540)$ | $4.698( \pm 0.540)$ | $0.646( \pm 0.040)$ | $M_{\odot}$ |
| $a_{3}\left(i^{\prime}=90^{\circ}\right)$ | $8.2( \pm 1.5)$ | $6.5( \pm 1.5)$ | $3.8( \pm 0.5)$ | au |
| $a_{3}\left(i^{\prime}=70^{\circ}\right)$ | $8.2( \pm 1.3)$ | $6.5( \pm 1.3)$ | $3.8( \pm 0.5)$ | au |
| $a_{3}\left(i^{\prime}=50^{\circ}\right)$ | $8.3( \pm 1.8)$ | $6.6( \pm 1.8)$ | $3.8( \pm 0.8)$ | au |
| $a_{3}\left(i^{\prime}=30^{\circ}\right)$ | $8.8( \pm 2.5)$ | $7.1( \pm 2.5)$ | $3.9( \pm 0.8)$ | au |
| $a_{3}\left(i^{\prime}=10^{\circ}\right)$ | 11.7( $\pm 3.2)$ | 10.3( $\pm 3.2)$ | $4.3( \pm 1.1)$ | au |

most plausible explanation for the cyclic period changes is the light travel time effect (LTTE) through the presence of a third body. By assuming a circular orbit, expressed by the equation
$f(m)=\frac{4 \pi^{2}}{G T_{3}^{2}} \times\left(a_{12}^{\prime} \sin i^{\prime}\right)^{3}$,
(where $a_{12}^{\prime} \sin i^{\prime}=A_{3} \times c$ and $c$ is the speed of light), the mass function for the assumed tertiary component is computed. The following equation is then used
$f(m)=\frac{\left(m_{3} \sin i^{\prime}\right)^{3}}{\left(M_{1}+M_{2}+m_{3}\right)^{2}}$
By taking into account the physical parameters, we obtained (section 3) the masses and the orbital radii of the third companion by computation. The values for several different orbital inclinations ( $i^{\prime}$ ) are given in table 4 . As shown in this table, the assumed tertiary component is invisible unless the orbital inclination, $i^{\prime}$, is very small $\left(i^{\prime}<10^{\circ}\right)$. If the tertiary companion is coplanar to the eclipsing pair (i.e., with the same inclination as the eclipsing binary), its mass should be $m_{3}=0.097 \pm 0.019 M_{\odot}$, which is too small to be detected.

### 4.2 AXCas

AX Cas has been monitored for a long time since 1937. We collected all of 60 times of minimum light, consisting of 14 visual and photographic data and 46 photoelectric and CCD ones (table 5). Because the visual and photographic data showed a very large scatter, we only used
the 46 CCD data points to correct the period with the ephemeris, $2456220.22457+0.600367560 \times E$, where $T_{0}$ is one of our minimum value, and $P_{0}$ was obtained from Alfonso-Garzón et al. (2012). The new linear ephemeris is

$$
\begin{align*}
\operatorname{Min} . \mathrm{I}(\mathrm{HJD})= & 2456220.2205( \pm 0.0010) \\
& +0.60037157( \pm 0.00000031) \times E \tag{6}
\end{align*}
$$

We then used this linear ephemeris to calculate the $(O-C)$ values for all 60 data points, yielding an $(O-C)$ diagram, as shown in figure 6 , where the solid cycles refer to the photoelectric or CCD primary minima, and the open ones refer to the photoelectric or CCD secondary minima; the crosses denote the visual or photographic minima, meanwhile; the solid line represents an eccentric ephemeris variation. A long-term period variation can be very clearly seen in figure 6. A parabola formula was used to fit this $(O-C)$ diagram, yielding a fitting of

Min. $I(H J D)=2456220.2204( \pm 0.0004)$

$$
\begin{align*}
& +0.600371245( \pm 0.00000003) \times E \\
& -5.462( \pm 0.007) \times 10^{-11} \times E^{2} \tag{7}
\end{align*}
$$

By using the quadratic term in this equation, a possible secular period increase rate is determined as being $d P / d t=$ $-6.646( \pm 0.009) \times 10^{-8} \mathrm{~d} \mathrm{yr}^{-1}$. However, we found that the residuals of the $(O-C)$ still showed a cyclic variation. Hence, we calculated the $(O-C)_{1}$ values for all 60 data first using equation (7), then performing an eccentric orbit fitting to simulate the photoelectric and CCD $(O-C)_{1}$

Table 5. Times of light minima for AX Cas.

| JD (Hel.) | Method | Min. | Epoch | $(\mathrm{O}-\mathrm{C})(\mathrm{d})$ | $(\mathrm{O}-\mathrm{C})_{1}(\mathrm{~d})$ | Residuals (d) | Reference |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2428626.43100 | photog. | prim. | -45961 | -0.11177 | -0.01116 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2428635.45000 | photog. | prim. | -45946 | -0.09834 | 0.00219 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2429114.53800 | photog. | prim. | -45148 | -0.10686 | -0.01003 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2429162.58600 | photog. | prim. | -45068 | -0.08858 | 0.00787 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2429165.59700 | photog. | prim. | -45063 | -0.07944 | 0.01699 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2429883.60800 | photog. | prim. | -43867 | -0.11284 | -0.02182 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2429903.46200 | photog. | prim. | -43834 | -0.07110 | 0.01976 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2429932.26400 | photog. | prim. | -43786 | -0.08694 | 0.00371 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2430100.36500 | photog. | prim. | -43506 | -0.08998 | -0.00056 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2430462.39200 | photog. | prim. | -42903 | -0.08703 | -0.00027 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2430839.43100 | photog. | prim. | -42275 | -0.08138 | 0.00266 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2431001.52600 | photog. | prim. | -42005 | -0.08670 | -0.00381 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2431145.60900 | photog. | prim. | -41765 | -0.09288 | -0.01101 |  | $(\mathrm{O}-\mathrm{C})$ gate way* |
| 2448562.47700 | vis | prim. | -12755 | -0.00412 | 0.00076 |  | BAV No. 60 |
| 2450752.62980 | CCD | prim. | -9107 | -0.00681 | -0.00509 | -0.00498 | BBSAG No. 116 |
| 2451336.20300 | CCD | prim. | -8135 | 0.00522 | 0.00633 | -0.00066 | ( $\mathrm{O}-\mathrm{C}$ ) gate way* |
| 2451430.46300 | CCD | prim. | -7978 | 0.00689 | 0.00791 | 0.00006 | IBVS No. 5016 |
| 2451780.48370 | CCD | prim. | -7395 | 0.01096 | 0.01168 | 0.00187 | IBVS No. 5296 |
| 2452576.57300 | CCD | prim. | -6069 | 0.00756 | 0.00773 | 0.00020 | IBVS No. 5378 |
| 2452992.62680 | CCD | prim. | -5376 | 0.00386 | 0.00383 | -0.00063 | IBVS No. 5502 |
| 2453215.36440 | CCD | prim. | -5005 | 0.00361 | 0.00348 | 0.00074 | IBVS No. 5731 |
| 2453254.38610 | CCD | prim. | -4940 | 0.00116 | 0.00102 | -0.00141 | IBVS No. 5657 |
| 2454085.29620 | CCD | prim. | -3556 | -0.00300 | -0.00332 | 0.00061 | IBVS No. 5657 |
| 2454092.49970 | CCD | prim. | -3544 | -0.00396 | -0.00428 | -0.00029 | IBVS No. 5657 |
| 2454097.30370 | CCD | prim. | -3536 | -0.00293 | -0.00325 | 0.00076 | IBVS No. 5781 |
| 2454367.46880 | CCD | prim. | -3086 | -0.00503 | -0.00537 | 0.00040 | IBVS No. 5830 |
| 2454388.48350 | CCD | prim. | -3051 | -0.00334 | -0.00368 | 0.00220 | IBVS No. 5830 |
| 2454390.28310 | CCD | prim. | -3048 | -0.00485 | -0.00519 | 0.00070 | IBVS No. 5830 |
| 2454405.29190 | CCD | prim. | -3023 | -0.00534 | -0.00568 | 0.00029 | IBVS No. 5889 |
| 2454704.87190 | CCD | prim. | -2524 | -0.01076 | -0.01109 | -0.00397 | IBVS No. 5875 |
| 2454752.30070 | CCD | prim. | -2445 | -0.01131 | -0.01164 | -0.00444 | IBVS No. 5889 |
| 2454776.31950 | CCD | prim. | -2405 | -0.00737 | -0.00770 | -0.00047 | IBVS No. 5889 |
| 2454847.46690 | CCD | sec. | -2286.5 | -0.00401 | -0.00432 | 0.00293 | IBVS No. 5918 |
| 2455154.55560 | CCD | prim. | -1775 | -0.00536 | -0.00562 | 0.00082 | IBVS No. 5941 |
| 2455473.35500 | CCD | prim. | -1244 | -0.00327 | -0.00344 | 0.00052 | IBVS No. 6010 |
| 2455479.36000 | CCD | prim. | -1234 | -0.00198 | -0.00216 | 0.00174 | IBVS No. 5984 |
| 2455482.36010 | CCD | prim. | -1229 | -0.00374 | -0.00391 | -0.00003 | B.R.N.O. No. 37 |
| 2455491.36780 | CCD | prim. | -1214 | -0.00161 | -0.00178 | 0.00200 | IBVS No. 5984 |
| 2455491.66050 | CCD | sec. | -1213.5 | -0.00910 | -0.00927 | -0.00548 | IBVS No. 5984 |
| 2455514.48530 | CCD | sec. | -1175.5 | 0.00158 | 0.00141 | 0.00495 | IBVS No. 5984 |
| 2455792.45730 | CCD | sec. | -712.5 | 0.00154 | 0.00147 | 0.00170 | B.R.N.O. No. 38 |
| 2455794.55580 | CCD | prim. | -709 | -0.00126 | -0.00132 | -0.00111 | B.R.N.O. No. 38 |
| 2455797.55730 | CCD | prim. | -704 | -0.00161 | -0.00167 | -0.00149 | B.R.N.O. No. 38 |
| 2455799.35880 | CCD | prim. | -701 | -0.00123 | -0.00129 | -0.00114 | B.R.N.O. No. 38 |
| 2455802.36020 | CCD | prim. | -696 | -0.00169 | -0.00174 | -0.00163 | B.R.N.O. No. 38 |
| 2455805.36250 | CCD | prim. | -691 | -0.00125 | -0.00130 | -0.00122 | B.R.N.O. No. 38 |
| 2455820.37150 | CCD | prim. | -666 | -0.00153 | -0.00158 | -0.00170 | IBVS No. 6026 |
| 2455830.28170 | CCD | sec. | -649.5 | 0.00253 | 0.00248 | 0.00222 | B.R.N.O. No. 38 |
| 2455831.48280 | CCD | sec. | -647.5 | 0.00289 | 0.00284 | 0.00257 | B.R.N.O. No. 38 |
| 2455835.38070 | CCD | prim. | -641 | -0.00162 | -0.00167 | -0.00198 | IBVS No. 6026 |
| 2455838.38330 | CCD | prim. | -636 | -0.00088 | -0.00092 | -0.00127 | B.R.N.O. No. 38 |
| 2455858.50250 | CCD | sec. | -602.5 | 0.00587 | 0.00583 | 0.00521 | IBVS No. 6026 |
| 2455875.31200 | CCD | sec. | -574.5 | 0.00497 | 0.00493 | 0.00409 | B.R.N.O. No. 38 |
| 2455905.62820 | CCD | prim. | -524 | 0.00240 | 0.00238 | 0.00114 | IBVS No. 6011 |
| 2456154.48730 | CCD | sec. | -109.5 | 0.00749 | 0.00758 | 0.00315 | B.R.N.O. No. 38 |

Table 5. (Continued)

| JD (Hel.) | Method | Min. | Epoch | $(O-C)(\mathrm{d})$ | $(O-C)_{1}(\mathrm{~d})$ | Residuals (d) |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2456179.39980 | CCD | prim. | -68 | 0.00457 | 0.00468 | -0.00004 | Reference |
| 2456214.22009 | CCD | prim. | -10 | 0.00331 | 0.00344 | -0.00169 |  |
| 2456220.22457 | CCD | prim. | 0 | 0.00407 | 0.00420 | -0.00100 | The paper |
| 2456292.26870 | CCD | prim. | 120 | 0.00361 | 0.00378 | -0.00224 | The paper |
| 2456535.42480 | CCD | prim. | 525 | 0.00923 | 0.00954 | JAAVSO No. 42 |  |

*〈http://var.astro.cz/ocgate/ocgate.php?star=ax+cas\&submit=Submit\&lang=en〉.


Fig. 6. $(O-C)$ diagram of $A X C$ Cas formed by all available measurements. The $(O-C)$ values were computed by using a newly determined linear ephemeris, equation (6). The solid circles refer to the photoelectric or CCD primary minima and the open ones to the photoelectric or CCD secondary minima; the crosses denote visual or photographic minima, while the solid line represents an eccentric ephemeris variation.


Fig. 7. $(O-C)_{1}$ diagram of AXCas determined by equation (7). The symbols are the same as those in figure 4.


Fig. 8. CCD residuals, equation (8), of $(O-C)_{1}$ diagram of $A X C$ Cas. The symbols are the same as those in figure 4.
values. The result is:

$$
\begin{align*}
(O-C)_{1}= & 0.0014( \pm 0.0004) \\
& +0.0031( \pm 0.0002) \cos 0.0428 E \\
& +0.0078( \pm 0.0018) \sin 0.0428 E \\
& +0.0007( \pm 0.0018) \cos 0.0856 E \\
& +0.0008( \pm 0.0008) \sin 0.0856 E \tag{8}
\end{align*}
$$

The corresponding $(O-C)_{1}$ curve is shown in figure 7, where the symbols are the same as those in figure 6. The corresponding residuals are plotted in figure 8. With the aid of the following relations (Kopal 1978):
$\omega=360^{\circ} P_{\mathrm{e}} / T$,
$A_{3}=\sqrt{a_{1}^{2}+b_{1}^{2}}$,
$e^{\prime}=2 \sqrt{\frac{a_{2}^{2}+b_{2}^{2}}{a_{1}^{2}+b_{1}^{2}}}$,
$w^{\prime}=\arctan \frac{\left(b_{1}^{2}-a_{1}^{2}\right) b_{2}+2 a_{1} b_{1} a_{2}}{\left(a_{1}^{2}-b_{1}^{2}\right) a_{2}+2 a_{1} b_{1} b_{2}}$,
the period of the orbital oscillation was determined to be $T_{3}=13.8 \mathrm{yr}$ with an eccentricity of $e^{\prime}=0.2533 ; P_{\mathrm{e}}$ is the ephemeris period ( 0.60037157 ); $e^{\prime}$ is the orbital eccentricity; $w^{\prime}$ is the longitude of the orbital periastron, and $a_{1}$, $b_{1}, a_{2}$, and $b_{2}$ are the corresponding fitting coefficients ( $a_{1}=$ $0.0031, b_{1}=0.0078, a_{2}=0.0007$, and $\left.b_{2}=0.0008\right)$. This eccentrical period oscillation may be caused by the light travel-time effect (LTTE) of a tertiary component. Then, by using equations (4) and (5), and taking into account the physical parameters $M_{1}=2.0 M_{\odot}$, and $M_{2}=0.45 M_{\odot}$ (M. A. Svechnikov \& Eh. F. Kuznetsova 1990), ${ }^{4}$ the masses and the orbital radii of the third companion were computed. The values for several different orbital inclinations ( $i^{\prime}$ ) are given in table 4. If the tertiary companion is coplanar to the eclipsing pair (i.e., with the same inclination as the eclipsing binary), its mass should be $m_{3}=0.54 \pm 0.03 M_{\odot}$.

However, the primary component of AX Cas may not be as large as $2.0 M_{\odot}$. As mentioned in section 1 , the temperature of the AX Cas' primary component would be $5600 \pm$ 370 K according to the colors, suggesting that $M_{1}=0.90 \pm$ $0.16 M_{\odot}, R_{1}=0.90 R_{\odot}$, and $L_{1}=0.719 L_{\odot}$. If we agreed with the mass ratio ( $q=M_{2} / M_{1}=0.225$ ) determined by M. A. Svechnikov and Eh. F. Kuznetsova (1990), ${ }^{4}$ we obtain $M_{2}=0.20 \pm 0.04 M_{\odot}$. The new set of masses for the components of AX Cas were used to compute the mass of the third body. One can find this new result in table 4 . In this case, the mass of the $m_{3}$ is $0.34 \pm 0.03 M_{\odot}$.

Like V1107 Cas, the mass-transfer time scale, $\tau$, and the thermal time scale, $t_{\mathrm{th}}$, of AX Cas were estimated. If $M_{1}$ $=2.0 M_{\odot}$, we obtain $\tau \sim 3.01 \times 10^{7} \mathrm{yr}$ and $t_{\mathrm{th}}=1.06 \times$ $10^{7} \mathrm{yr}$; however, $M_{1}=0.90 M_{\odot}$, we have $\tau \sim 1.35 \times 10^{7} \mathrm{yr}$ and $t_{\mathrm{th}}=2.50 \times 10^{7} \mathrm{yr}$. This indicates that $\tau$ and $t_{\mathrm{th}}$ are of the same order.

## 5 Discussion and conclusions

The $V R_{\mathrm{c}} I_{\mathrm{c}}$ light-curve solution shows that the mass ratio, $q$, of V1107 Cas is $1.797 \pm 0.006$, while the fill-out factor is $15.2 \% \pm 1.8 \%$. The results suggest that V1107 Cas is a W-type moderate mass-ratio shallow contact binary. [The W-type is a contact binary in which the more massive component is the cooler one. Please see the classification by Binnendijk (1970)]. From the temperature ( $T_{2}$ ) of the more massive component, we estimated its mass to be $M_{2}=0.70 \pm 0.03$. The rest physical parameters are obtained by using the LC program of the W-D code (Wilson \& Devinney 1971; Wilson 1979, 1990, 1994, 2008; Van Hamme \& Wilson 2007). Thus, all of the parameters are: $M_{1}=0.39 \pm 0.01 M_{\odot}, M_{2}=0.70 \pm 0.03 M_{\odot}$, $R_{1}=0.52 \pm 0.10 R_{\odot}, R_{2}=0.68 \pm 0.12 R_{\odot}, L_{1}=0.178$ $\pm 0.108 L_{\odot}$, and $L_{2}=0.196 \pm 0.116 L_{\odot}$, where the errors
mainly come from the uncertainties in the temperatures of the components. Someone may ask, why you did not use the higher temperature ( $T_{1}$ ) of 5200 K to estimate $M_{1}$. We can explain this. The temperature of the less-massive component is affected by the more-massive one because of the mass transfer. The material flow heats the less-massive component so that it is much hotter than it would be ordinarily. Such an affected temperature is not suitable for estimating the mass value.

Based on all of the available eclipse times, the period variations of V1107 Cas and AXCas were investigated. All of them show long-term period decreases of $d P / d t=$ $-1.52( \pm 0.55) \times 10^{-7} \mathrm{~d} \mathrm{yr}^{-1}$ for V1107 Cas and $d P / d t=$ $-6.646( \pm 0.009) \times 10^{-8} \mathrm{~d} \mathrm{yr}^{-1}$ for AX Cas. Besides the long-term period decrease, their periods both indicate cyclic oscillations, with period $T_{3}$ of 6.74 years for V1107 Cas and that of 13.8 years for AX Cas. Especially, the oscillation of AX Cas' period is an eccentric one with an $e^{\prime}$ of 0.2533 . The long-term period decreases can be explained as being mass transfer from the more-massive component to the less-massive one, while the cyclic period oscillations can be interpreted as being the light travel-time effect (LTTE) caused by a tertiary companion. The masses and distance to the center binary are listed in table 4.

Based on the coplanar assumption, the $m_{3}$ in V1107 Cas system should be $0.097 \pm 0.019 M_{\odot}$, while in the AX Cas system it should be $0.34 \pm 0.03 M_{\odot}$. A main-sequence star with $0.097 \pm 0.019 M_{\odot}$ should be of the M7 type (Cox 2000). The luminosity of M5 is $0.001 L_{\odot}$, so the third body in V1107 Cas must be fainter than that. We can safely assume $l_{3}$ is $0.0005 L_{\odot}$. Such an $l_{3}$ is only a $0.13 \%$ contribution to the system, which agrees with the $0.16 \%$ reduced from the DC program.

Samec and his colleagues have studied several precontact binary systems (i.e., V2421 Cyg, Samec et al. 2014; ZZ Eri, Samec et al. 2015a; V1001 Cas, Samec et al. 2015b). Compared with their work, we believe that AX Cas is also a precontact binary, like V2421 Cyg (Samec et al. 2014). This is also supported by a decreased period rate. Thus, AX Cas is a precontact binary, while V1107 Cas is a shallow contact binary; they have adjoining evolutional stages. That makes the situation much more significance than if they are just in the same field.

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## References

Alfonso-Garzón, J., Domingo, A., Mas-Hesse, J. M., \& Giménez, A. 2012, A\&A, 548, A79
Applegate, J. H. 1992, ApJ, 385, 621
Beljawsky, S. 1931, Astron. Nachr., 243, 115
Binnendijk, L. 1970, Vistas Astron., 12, 217
Claret, A., \& Gimenez, A. 1990, A\&A, 230, 412
Cox, A. N. ed. 2000, Allen's Astrophysical Quantities, 4th ed. (New York: Springer)
Hoffman, D. I., Harrison, T. E., \& McNamara, B. J. 2009, AJ, 138, 466
Kopal, Z. 1978, Dynamics of Close Binary Systems (Dordrecht: D. Reidel Publishing Co.)

Lanza, A. F., Rodono, M., \& Rosner, R. 1998, MNRAS, 296, 893
Liao, W.-P., \& Qian, S.-B. 2010, MNRAS, 405, 1930
Lucy, L. B. 1967, Z. Astrophysik 65, 89
Ruciński, S. M. 1969, Acta Astron., 19, 245

Samec, R. G., Carrigan, B. J., Gray, J. D., French, J. A., McDermith, R. J., \& Padgen, E. E. 1998, AJ, 116, 895

Samec, R. G., Clark, J. D., Van Hamme, W., \& Faulkner, D. R. 2015a, AJ, 149, 48
Samec, R. G., Koenke, S. S., \& Faulkner, D. R. 2015b, AJ, 149, 30
Samec, R. G., Shebs, T. S., Faulkner, D. R., Van Hamme, W., \& Mathis, R. F. 2014, AJ, 147, 3
Tout, C. A., \& Hall, D. S. 1991, MNRAS, 253, 9
Van Hamme, W., \& Wilson, R. E. 2007, ApJ, 661, 1129
Wilson, R. E. 1979, ApJ, 234, 1054
Wilson, R. E. 1990, ApJ, 356, 613
Wilson, R. E. 1994, PASP, 106, 921
Wilson, R. E. 2008, ApJ, 672, 575
Wilson, R. E., \& Devinney, E. J. 1971, ApJ, 166, 605
Worthey, G., \& Lee, H.-c. 2011, ApJS, 193, 1
Xiang, F.-Y., Xiao, T.-Y., Zhang, B., \& Shi, X.-D. 2015a, AJ, 150, 9
Xiang, F.-Y., Xiao, T.-Y., \& Yu, Y.-X. 2015b, AJ, 150, 25
Xiang, F.-Y., Yu, Y.-X., \& Xiao, T.-Y. 2015c, AJ, 149, 62
Zhou, A.-Y., Jiang, X.-J., Zhang, Y.-P., \& Wei, J.-Y. 2009, Res. Astron. Astrophys., 9, 349
Zhu, L., \& Qian, S. 2006, MNRAS, 367, 423
Zhu, L., Qian, S.-B., Mikulǎšek, Z., Zejda, M., Zvěřina, P., \& Diethelm, R. 2010, AJ, 140, 215
Zhu, L.-Y., Qian, S.-B., \& Yang, Y.-G. 2008, AJ, 136, 337


[^0]:    ${ }^{2}$ Cutri，R．M．，et al．2003，2MASS All－Sky Catalog of Point Sources，VizieR Online Data Catalog，II／246 〈http：／／cdsarc．u－strasbg．fr／viz－bin／Cat？II／246〉．
    ${ }^{3}$ Zacharias，N．，Monet，D．G．，Levine，S．E．，Urban，S．E．，Gaume，R．，\＆Wycoff， G．L．Naval Observatory Merged Astrometric Dataset（NOMAD），VizieR Online Data Catalog，I／297 〈http：／／cdsarc．u－strasbg．fr／viz－bin／Cat？l／297〉．
    ${ }^{4}$ Svechnikov，M．A．，\＆Kuznetsova Eh．F．1990，Catalogue of approx－ imate photometric and absolute elements of eclipsing variable stars〈http：／／cdsarc．u－strasbg．fr／viz－bin／Cat？V／124〉．

[^1]:    ${ }^{5}$ Samus, N. N., et al. 2009, General Catalogue of Variable Stars, VizieR Online Data Catalog, B/gcvs 〈http://cdsarc.u-strasbg.fr/viz-bin/Cat?B/gcvs〉.

